#### Lecture 6

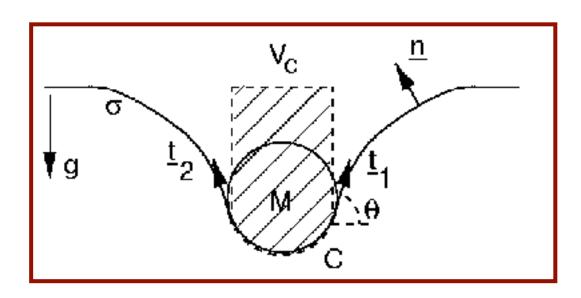
Floating bodies, capillary forces, rotating drops

### On floating bodies

Heavy things sink, light things float.

Not exactly.....

## Statics of floating bodies



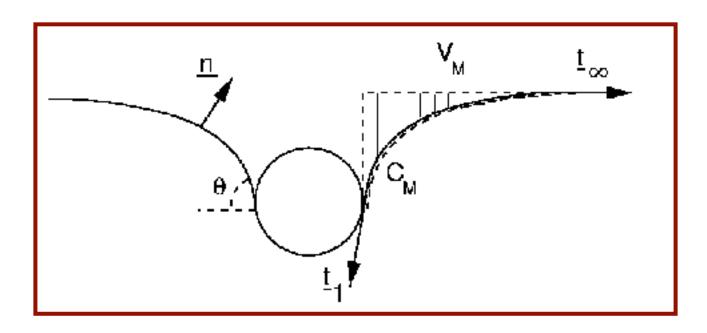
#### Force balance on body:

$$Mg = \hat{\mathbf{z}} \cdot \int_C -p \, \mathbf{n} \, d\ell + 2\sigma \sin \theta$$

**Buoyancy:** 
$$F_b = \int_C \rho g z (\hat{\mathbf{n}} \cdot \hat{\mathbf{z}}) d\ell = \rho g V_c$$

Surface tension:  $F_C = 2\sigma \sin \theta$ 

#### Force balance on meniscus



$$0 = \int_{C_{M}} -p \, \hat{\mathbf{n}} \cdot \hat{\mathbf{z}} \, d\ell = F_b + F_c$$

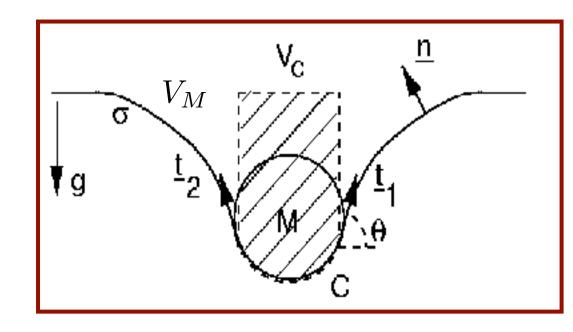
where 
$$F_b = \int_{C_M} \rho \, g \, z \, \hat{\mathbf{n}} \cdot \hat{\mathbf{z}} \, d\ell = \rho \, g \, V_M$$

$$F_C = \int_{C_M} \sigma \, (\nabla \cdot \hat{\mathbf{n}}) \, (\hat{\mathbf{n}} \cdot \hat{\mathbf{z}}) \, d\ell = 2\sigma (\hat{\mathbf{t}}_{\infty} - \hat{\mathbf{t}}_1) \cdot \hat{\mathbf{z}} = -2\sigma \sin \theta$$
via Frenet-Serret

### Generalized Archimedes Principle

 the weight of a floating body still equals that of the displaced fluid

$$Mg = F_B + F_C$$



Buoyancy: 
$$F_B = \rho g V_C$$

= wt of fluid displaced above body

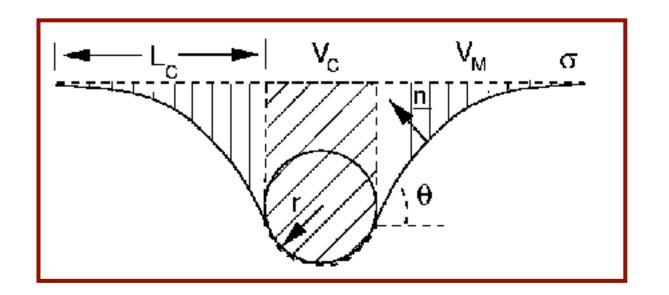
$$F_C = 2\sigma \sin \theta = \rho g V_M$$

= wt of fluid displaced above meniscus

### **Weight support:**

statics of floating bodies

- J. Keller (1998)



$$\Rightarrow F_b = \rho g V_c = \text{wt. of fluid displaced above body}$$

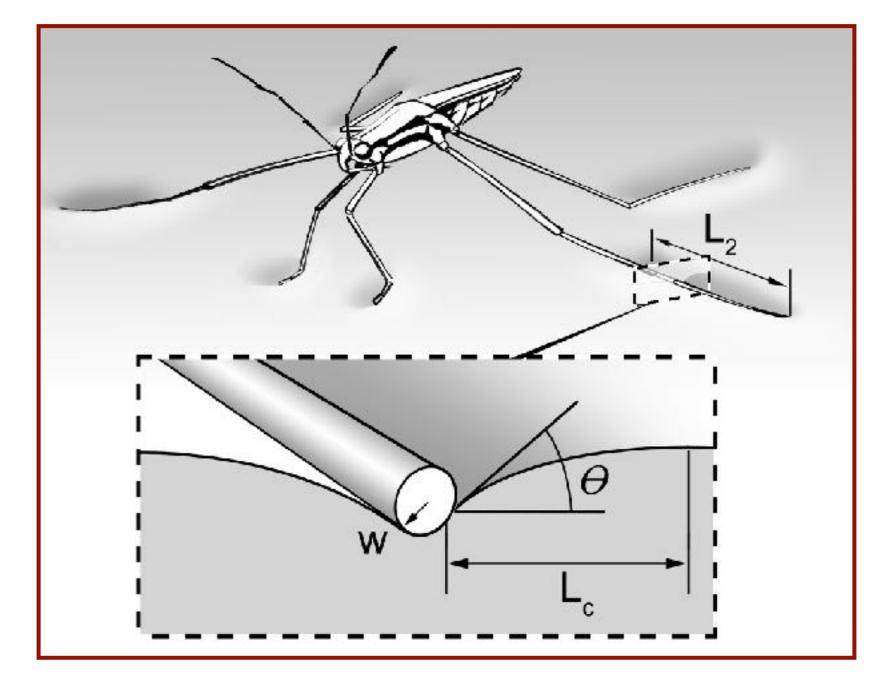
$$\Rightarrow F_c = 2\sigma \sin\theta = \rho g V_M = \text{ wt. of fluid above meniscus}$$

$$\implies \frac{F_b}{F_c} = \frac{V_c}{V_M} \approx \frac{r}{L_c} \qquad \text{where} \qquad L_c = \left(\frac{\sigma}{\rho g}\right)^{1/2} \approx 0.3 \text{ cm}$$

 $\Longrightarrow$  small creatures (eg. insects) supported principally by  $\sigma$ 

Capillary forces support the weight of water-walking insects.





Static weight support requires:  $Mg < 2\sigma P \sin \theta$ 

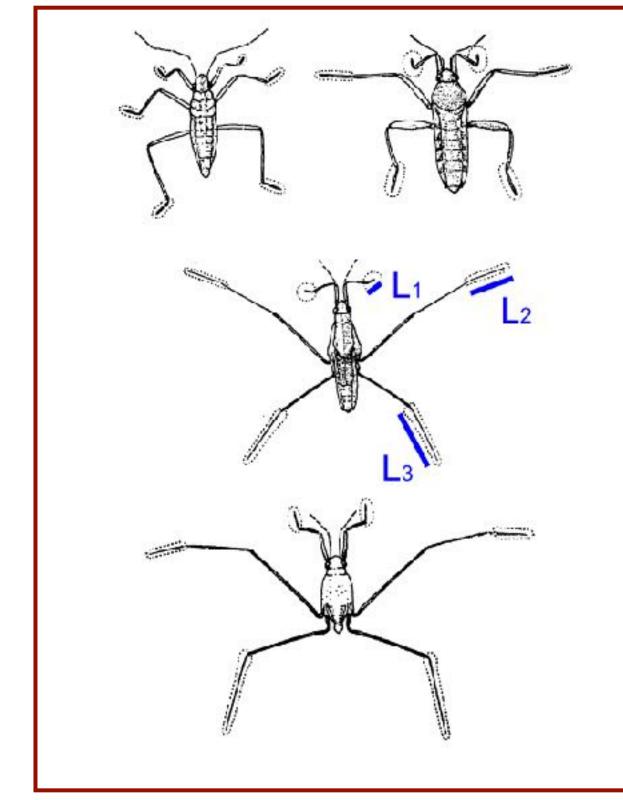
### Water striders

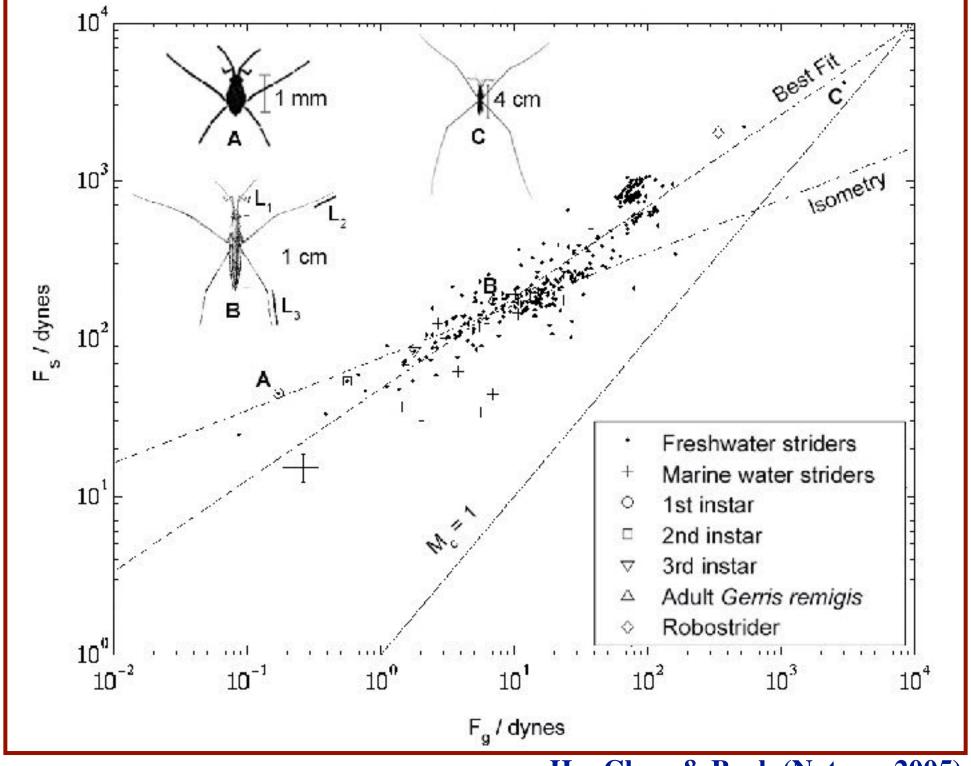
$$F_s = 2\sigma P$$
$$F_g = Mg$$

$$P = 2(L_1 + L_2 + L_3)$$

What is  $F_s(F_g)$ ?

i.e. what is the dependence of form on size?



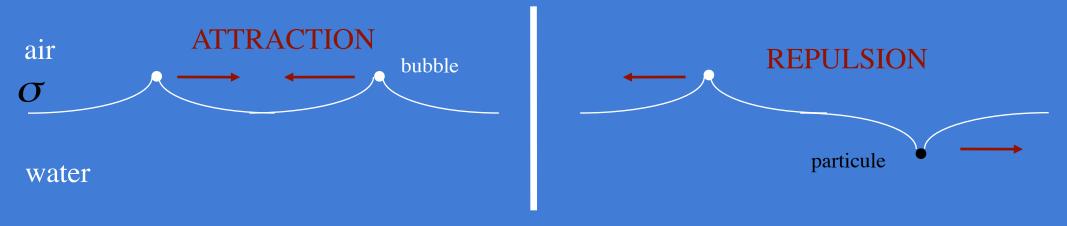


Hu, Chan & Bush (Nature, 2005)

# **Capillary forces**

### Capillary forces

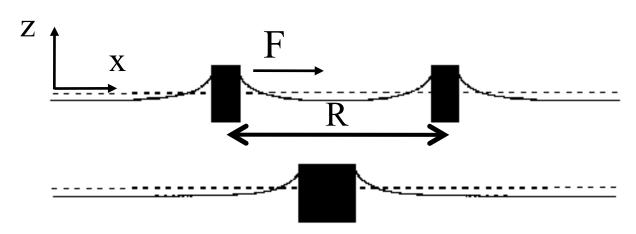
- exist by virtue of the interaction of the menisci of floating bodies
- attractive/repulsive if the menisci are of the same/opposite sense



- explains the formation of bubble rafts on champagne
- explains the mutual attraction of Cheerios, and their attraction to the walls
- utilized for self-assembly on the microscale

# Capillary attraction

$$E_{Tot} = E_S + GPE$$

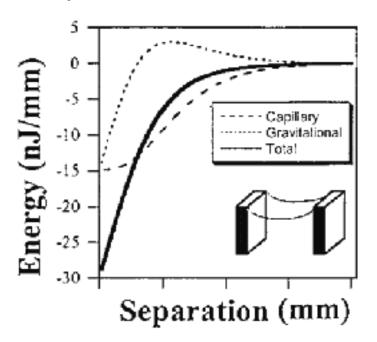


$$E_S = \sigma A(R)$$

$$GPE(R) = \int_{x=-\infty}^{x=\infty} \int_{z=0}^{h(x)} \rho g \, dz \, dx$$

$$\implies F(R) = -\frac{dE_{Tot}(R)}{dR}$$

#### Gryzbowski et al 2001



### Floating copper



How does it float? Why the attractive force between floaters?

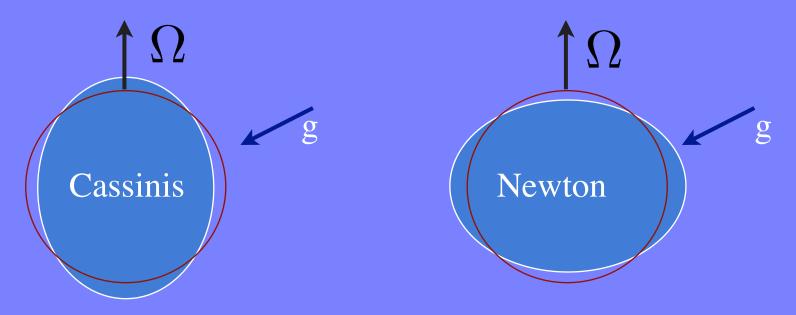
# The courtship of the Anurida



## On rotating fluid bodies

**Plateau droplets** 

#### The shape of a rotating, self-gravitating body



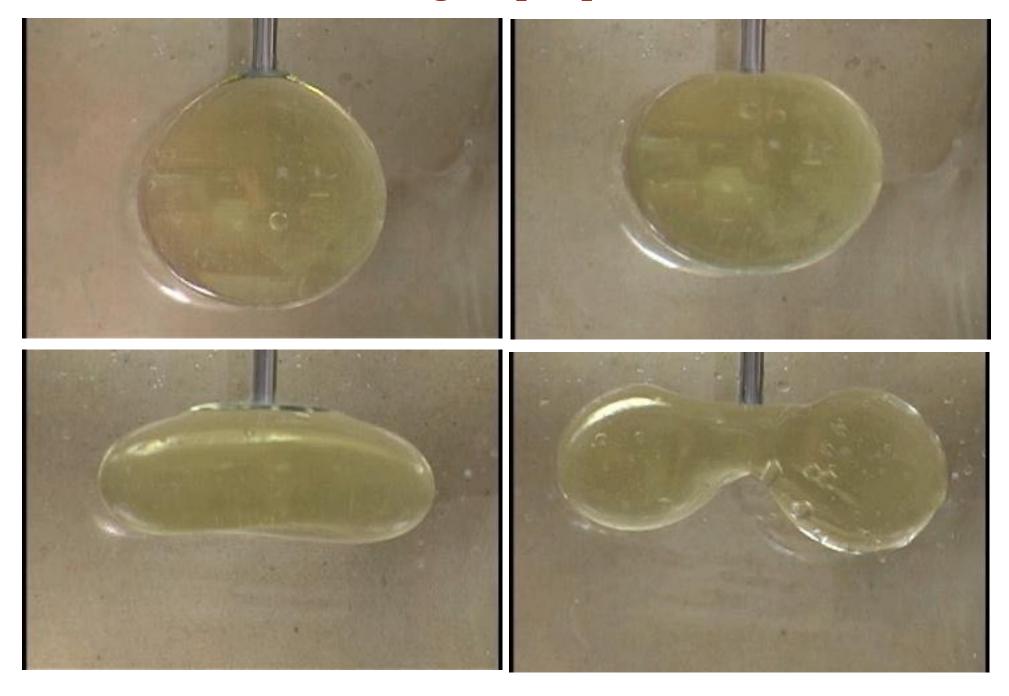
• debate was settled by geodetic measurements made in Lapland by Maupertuis, who was celebrated by Voltaire for having

"hammered down the poles and the Cassinis"

• later, Voltaire and Maupertuis had a falling out (over a woman, Emilie du Chatelet), and the former taunted the latter:

"You have gone to the ends of the Earth to confirm What Newton knew without leaving his home."

### Plateau's rotating drop experiments (1863)



The confining role of self-gravitation is played by surface tension.

**Rotating drops:** a model of celestial bodies

Normal force balance

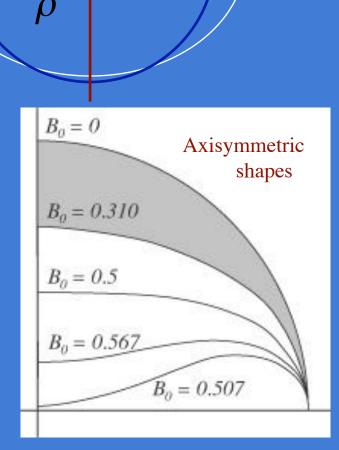
Plateau (1863) Chandrasekhar (1965)

$$\Delta \mathbf{p} + \frac{1}{2} \rho \Omega^2 r^2 = \sigma \nabla \cdot \mathbf{n}$$

Define: 
$$f(r,\theta) = z - h(r)$$

so that 
$$\mathbf{n} = \frac{\nabla f}{|\nabla f|} = \frac{\hat{\mathbf{z}} - h_r(r)\hat{\mathbf{r}}}{[1 + h_r^2(r)]^{1/2}} \quad \nabla \cdot \mathbf{n} = \frac{-rh_r - r^2h_{rr}}{r^2(1 + h_r^2)^{3/2}}$$

where 
$$\Delta P = a\Delta p/\sigma$$
 ,  $B_0 = \frac{\rho\Omega^2 a^3}{8\sigma} = \frac{\text{centrifugal}}{\text{curvature}}$ 



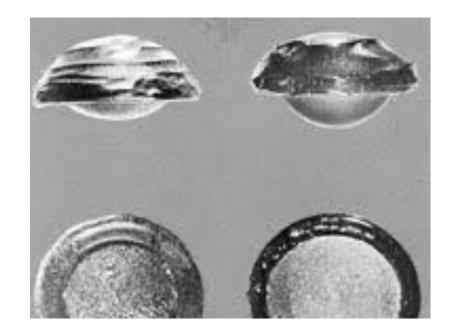
z = h(r)

- for 0.09 < B < 0.31, axisymmetric and lobed solutions possible
- for B > 0.31, only lobed forms obtain (Brown & Scriven, 1980)

#### **Tektites**

• centimetric rock forms produced by splash from ancient meteor impacts

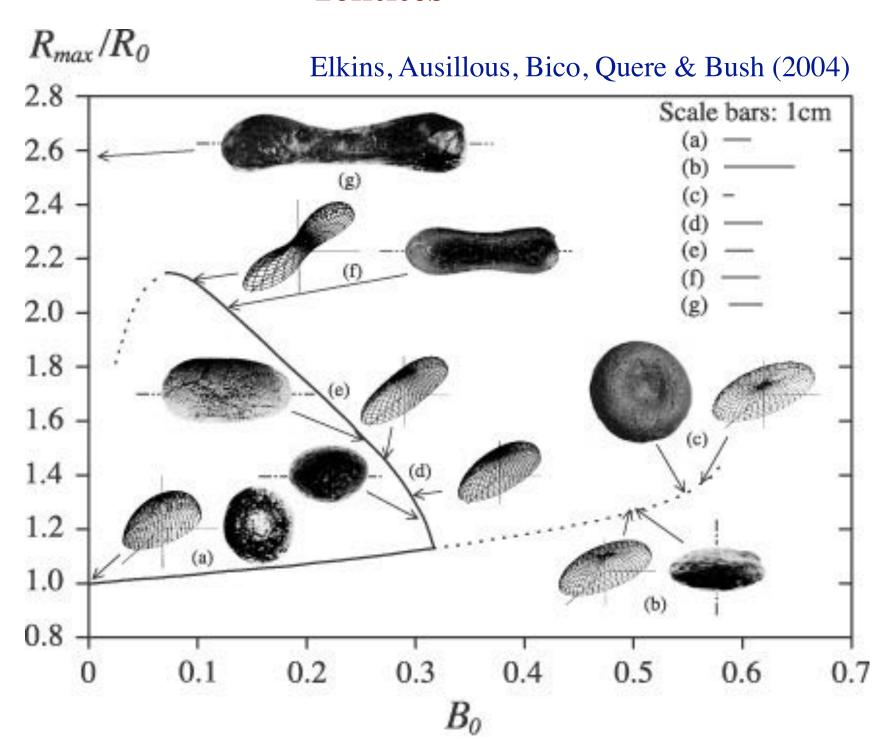




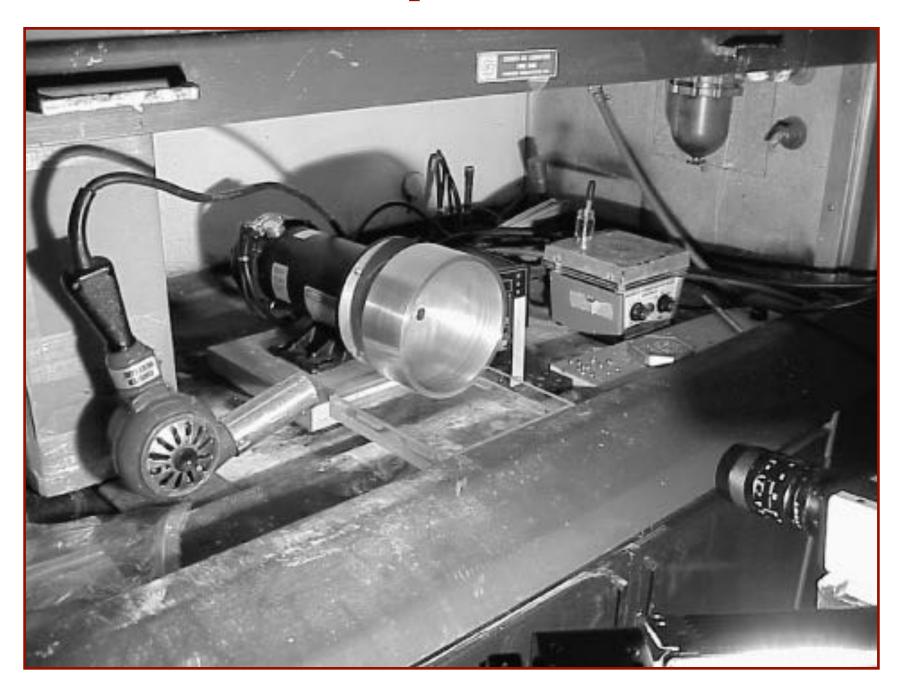




#### **Tektites**

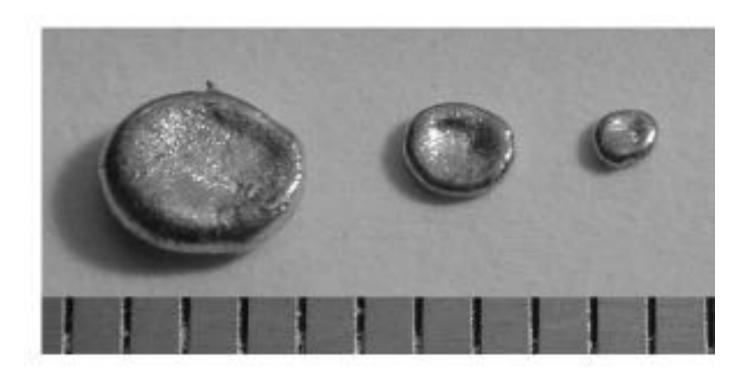


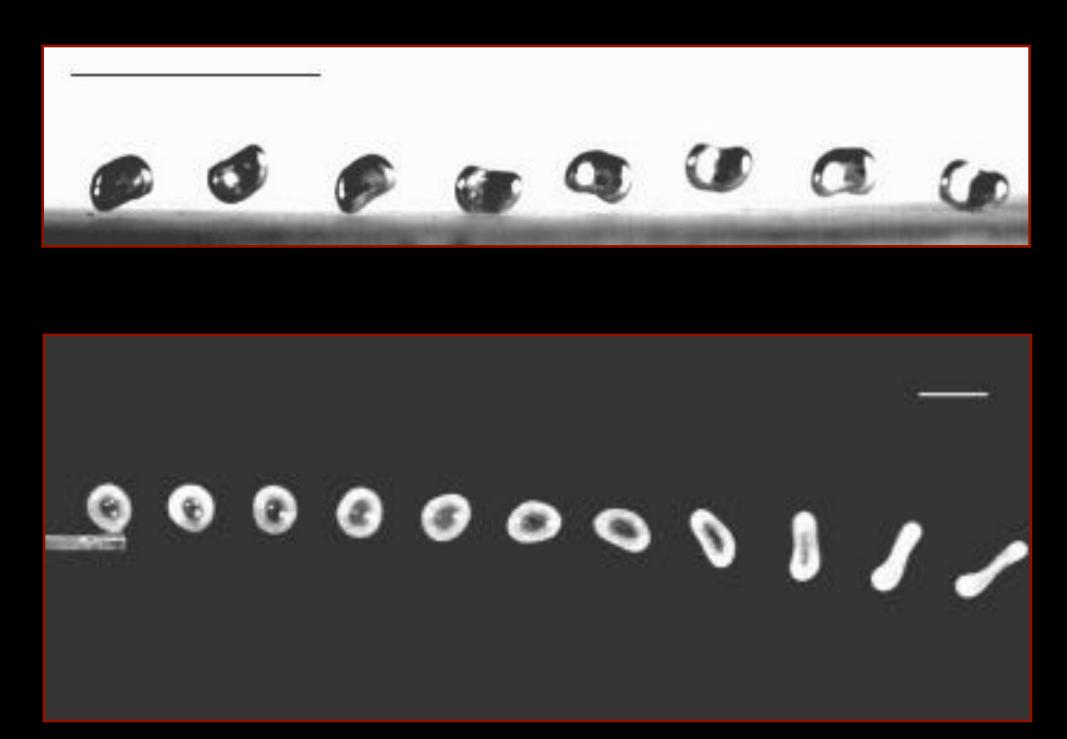
# A lab model of splash-form tektites



# Synthetic tektites







The tumbling tektite



The spinning drop tensiometer

