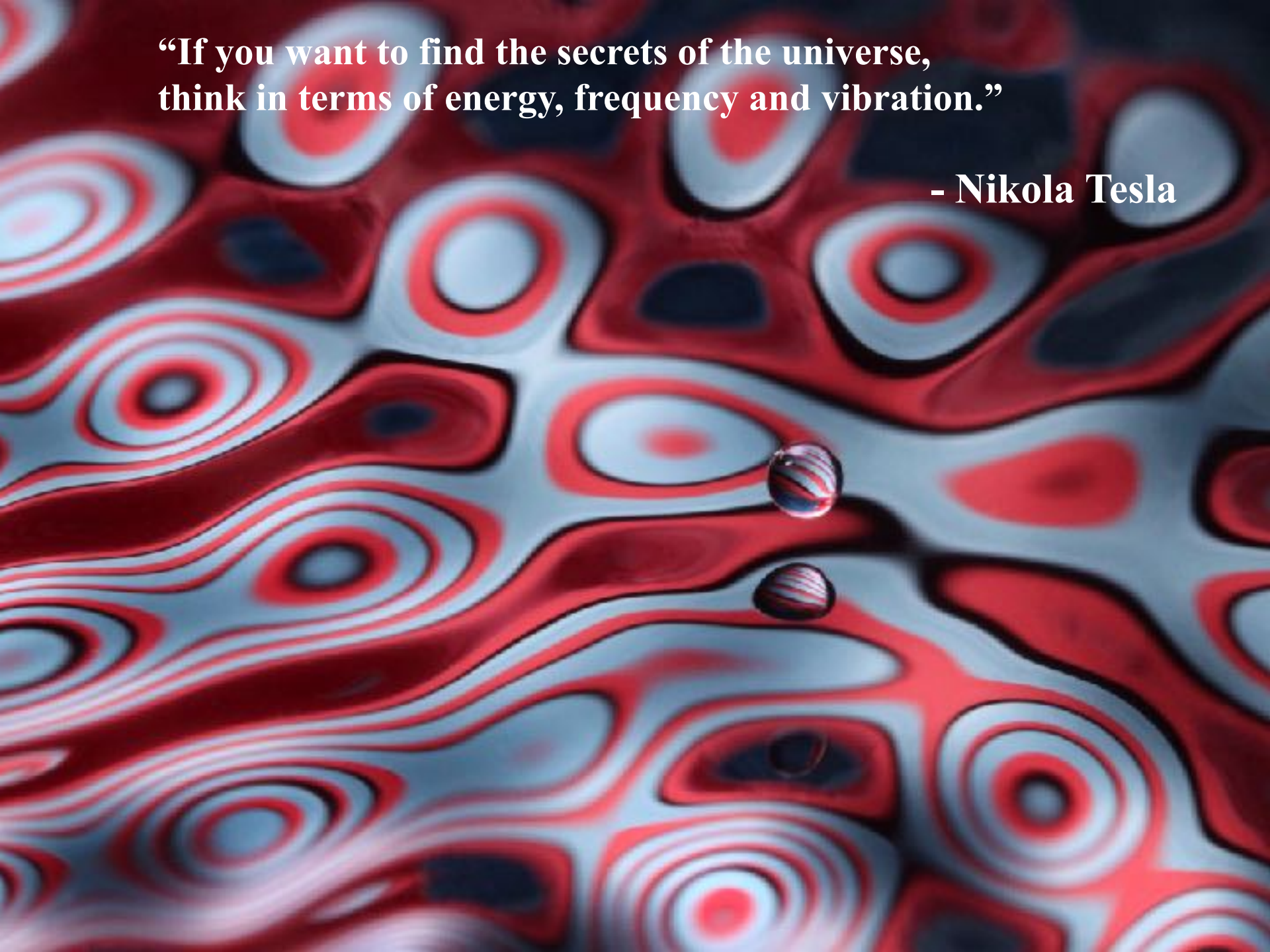


18.S996 Hydrodynamic quantum analogs

Lecture 11: Finale: the theory of walking droplets

**“If you want to find the secrets of the universe,
think in terms of energy, frequency and vibration.”**

- Nikola Tesla



Molacek: model summary

Vertical dynamics

$$m_D \ddot{Z} = -m_D g (1 + \Gamma \sin \omega t) \quad \text{when } Z \geq 0$$

$$m_D \ddot{Z} \left(1 + \frac{c_3}{Q^2(Z)} \right) + \frac{\mu R_0 c_2}{Q(Z)} \dot{Z} + \frac{3\sigma Z}{2Q(Z)} = -m_D g (1 + \Gamma \sin \omega t) \quad \text{otherwise}$$

Standing wave evolution

$$h(X, \tau) = \sum_{n=1}^N h_0(X, X_n, \tau, \tau_n)$$

$$h_0(X, X_n, \tau, \tau_n) \approx \frac{4\sqrt{2\pi} k_C^2 k_F \Omega h_e^{1/2}}{3 \cdot 3k_F^2 |E_0|} \left[\int F(u) \sin \frac{\Omega u}{2} du \right] \frac{H(\tau)}{\sqrt{\tau - \tau_n}} \exp \left\{ (1/\Gamma^p - 1) \frac{\tau - \tau_n}{\tau_d} \right\} J_0(k_C(X - X_n))$$

$$X_n = \int_{\tau_e} F(u) X(u) du / \int_{\tau_e} F(u) du, \quad \tau_n = \int_{\tau_e} F(u) u du / \int_{\tau_e} F(u) du$$

Horizontal dynamics

$$m_D \ddot{X} + \left[C F_N \sqrt{\rho R_0 / \sigma} + 6\pi \mu_a R_0 \right] \dot{X} = - \frac{\partial h}{\partial X} F_N$$

contact drag

air drag

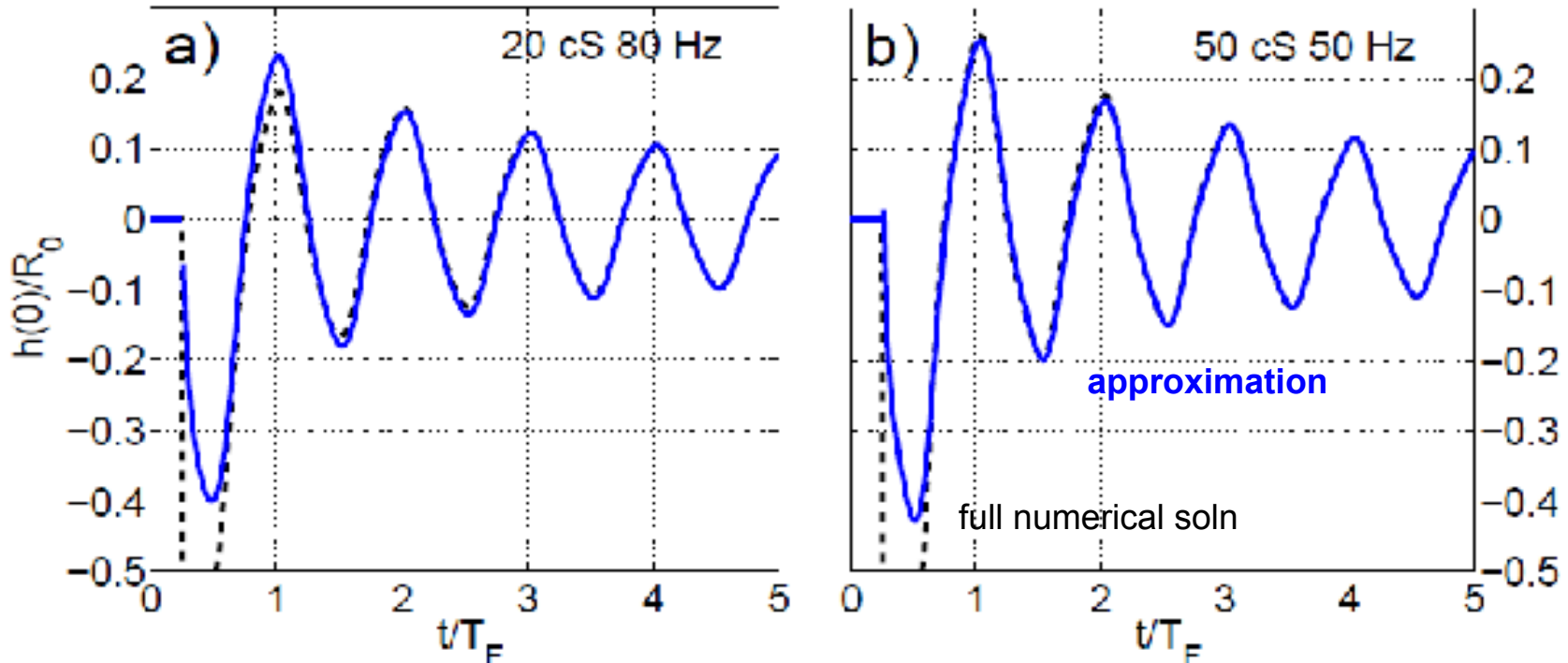
sloped
interface

Standing wave field

- standing wave created by a single drop impact:

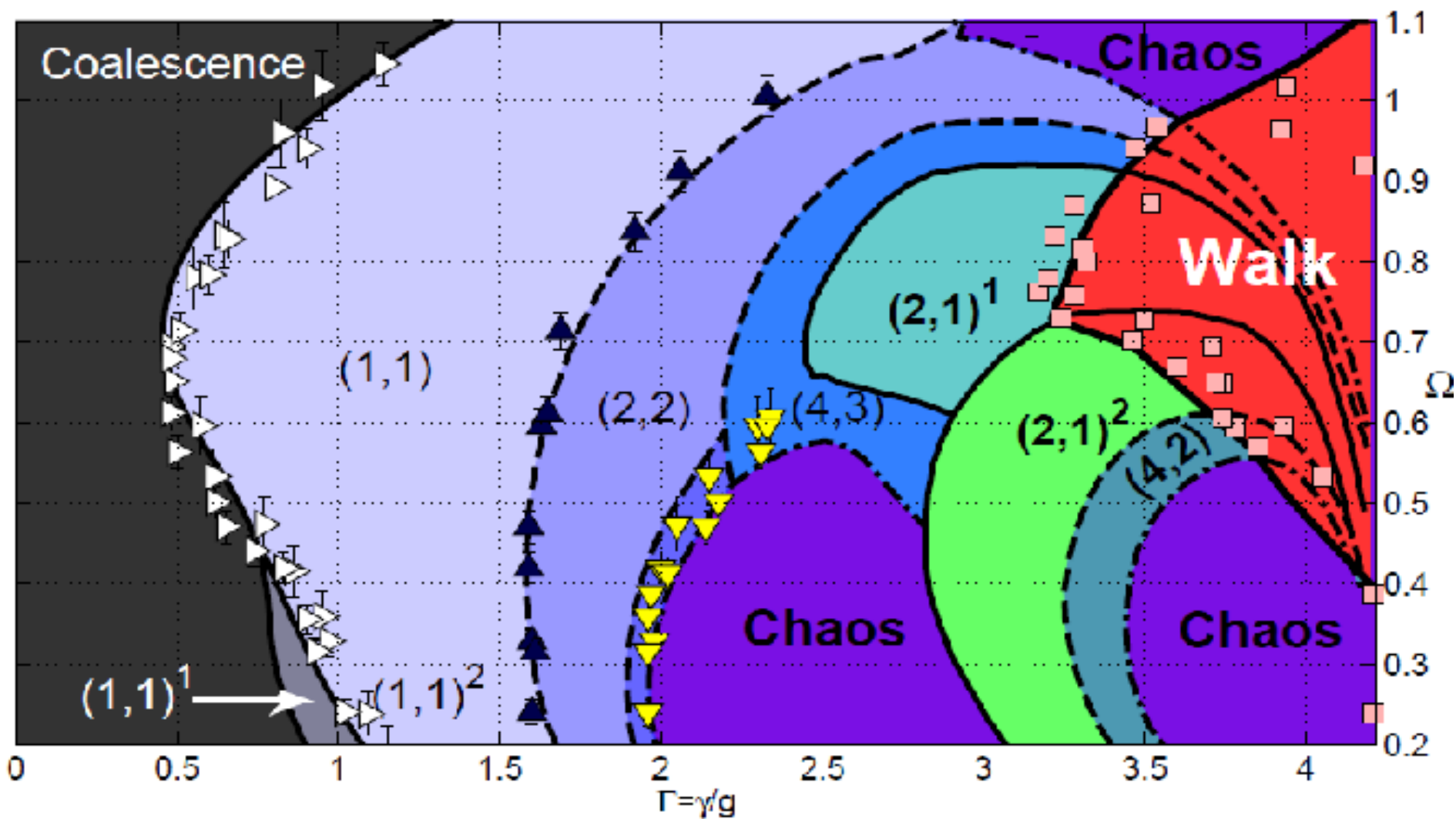
$$h(r, \tau) \approx \frac{1\sqrt{2\pi} k_C^2 k_F \Omega h_e^{1/2}}{3\sqrt{\tau} (3k_F^2 + \mathbb{B}\Omega)} \left[\int F(u) \sin \frac{\Omega u}{2} du \right] \cos \frac{\Omega \tau}{2} \exp \left\{ \left(\frac{\Gamma}{\Gamma_F} - 1 \right) \frac{\tau}{\tau_d} \right\} J_0(k_C r) .$$

- most unstable mode $k_C \approx k_F$ where $k_C^3 + \mathbb{B}\Omega \cdot k_F = \frac{1}{4}\Omega^2$

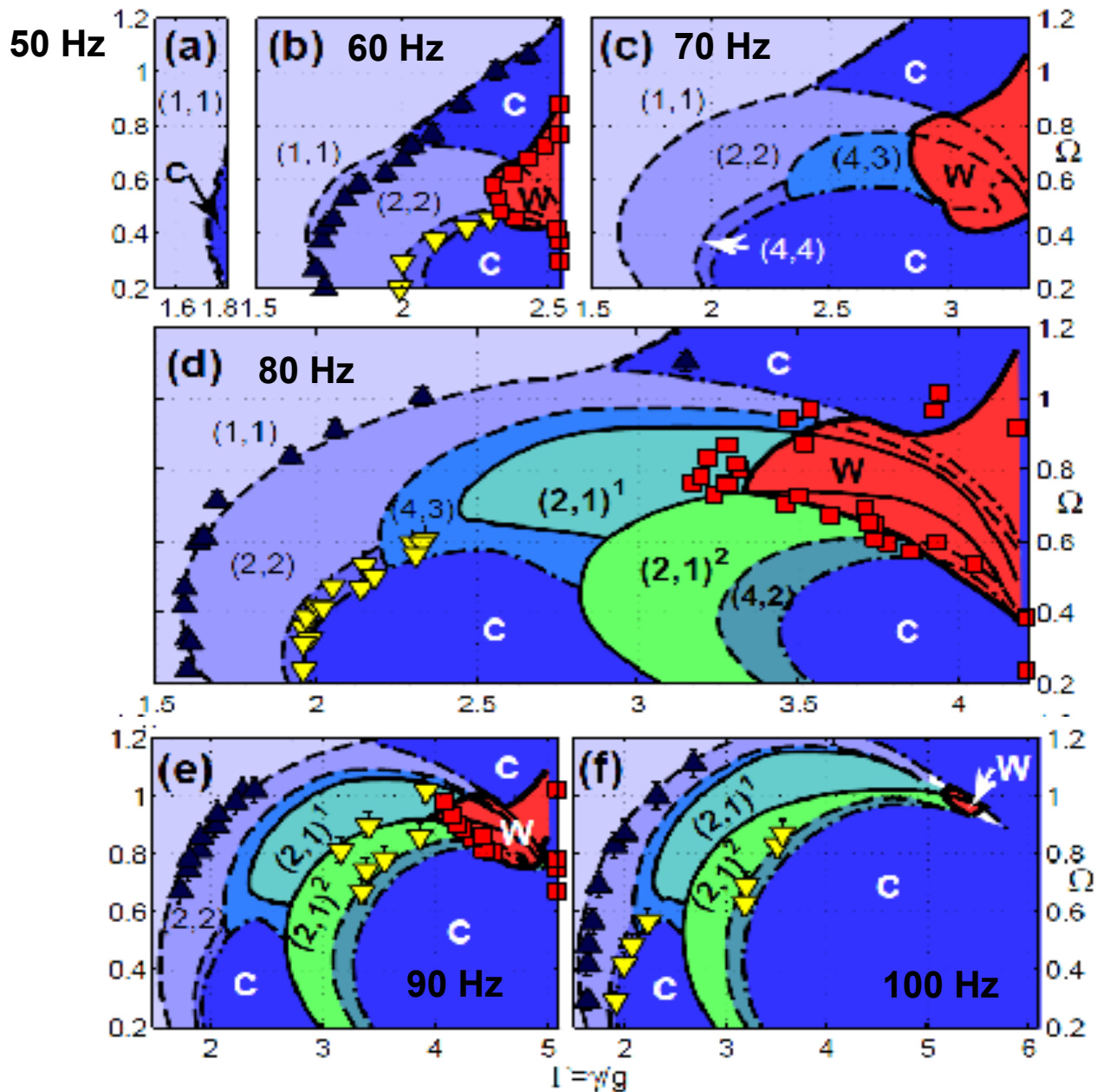


Predicted regime diagrams

20 cS, 80 Hz

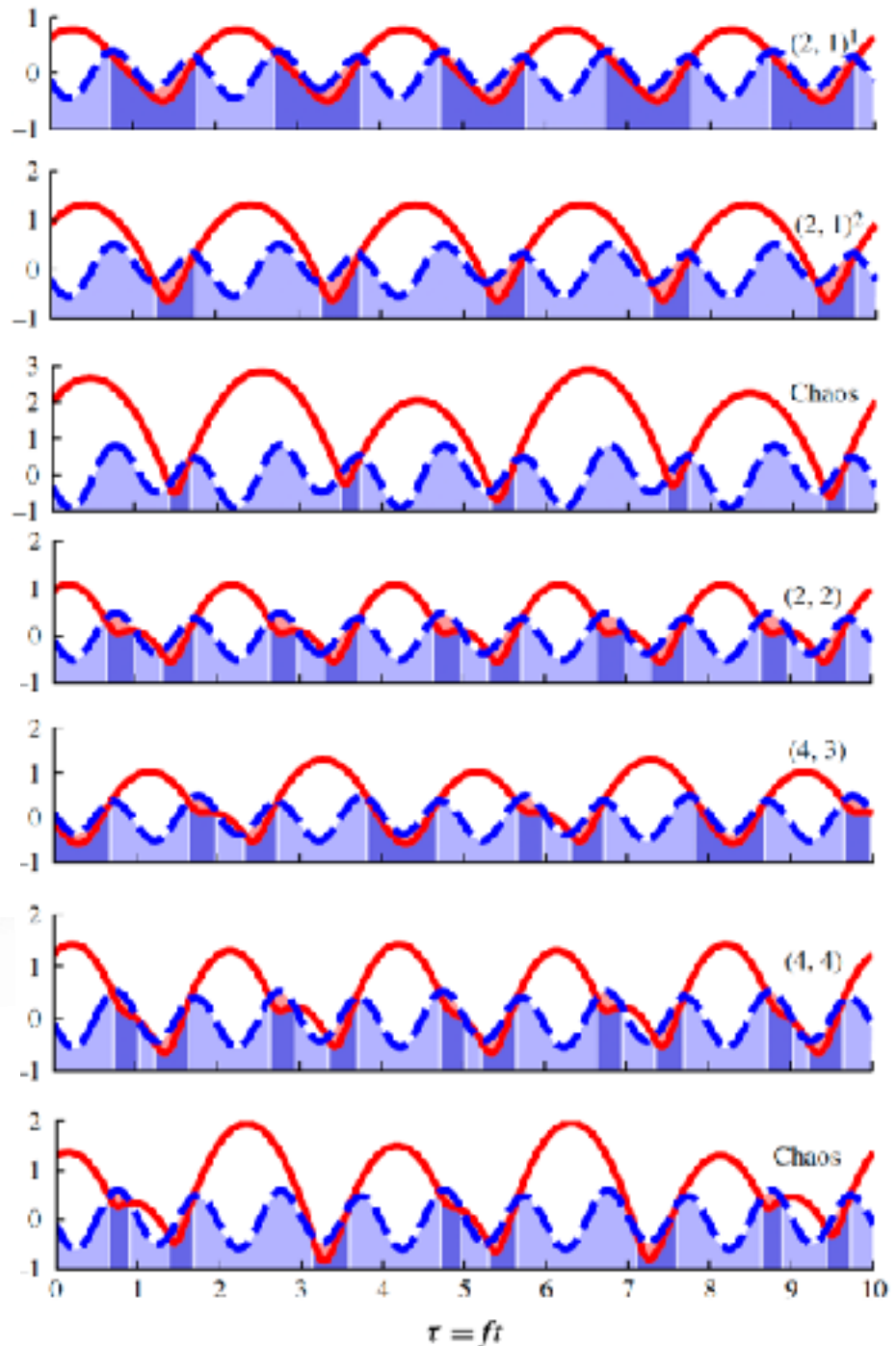


Predicted regime diagrams for 20 cS

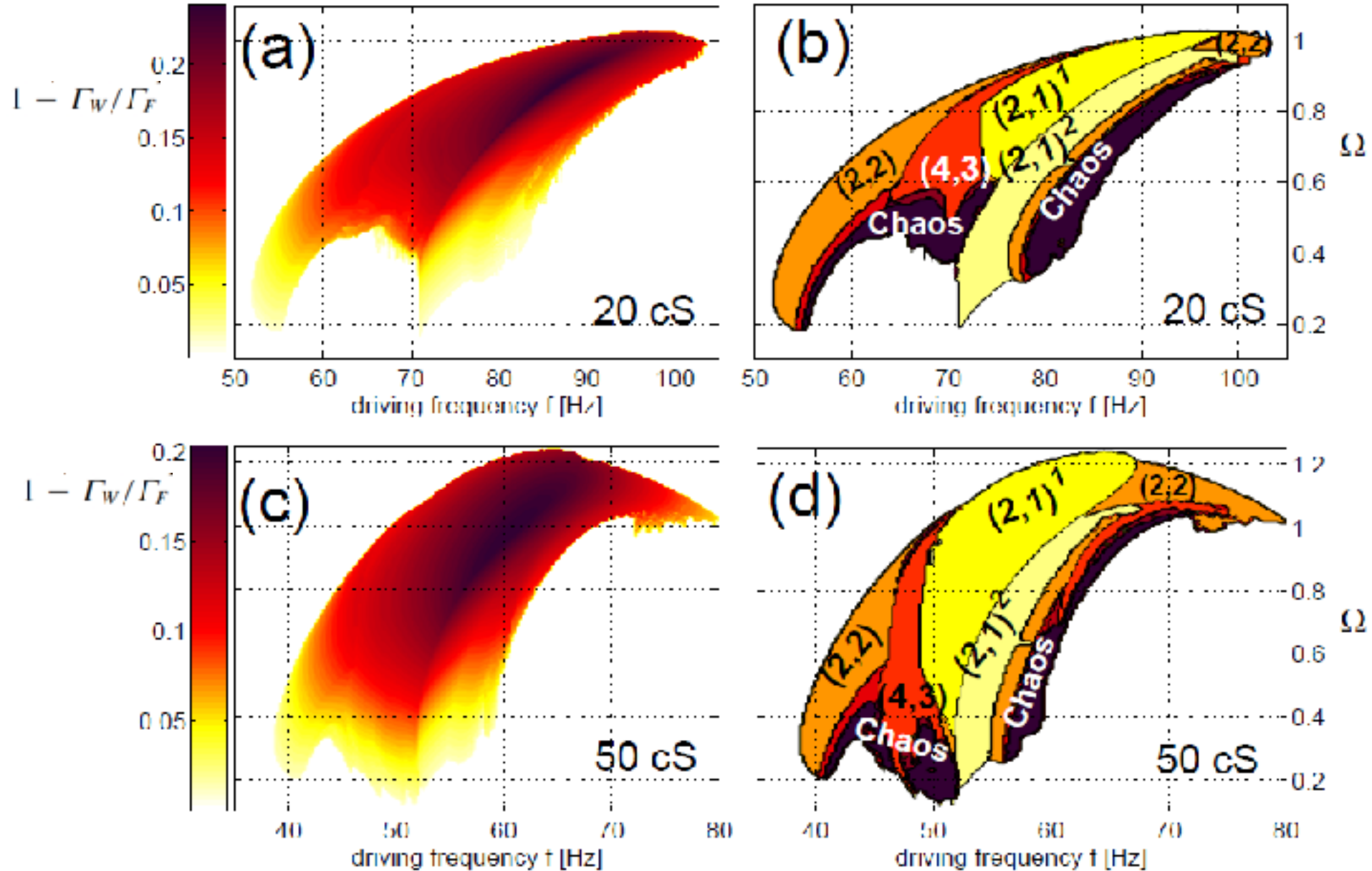


Walking modes

arising at various f for 20 cS



The walking regime



PHYSICS OF FLUIDS 25, 082002 (2013)

Exotic states of bouncing and walking droplets

Øistein Wind-Willassen,¹ Jan Moláček,² Daniel M. Harris,²
and John W. M. Bush^{2,a)}

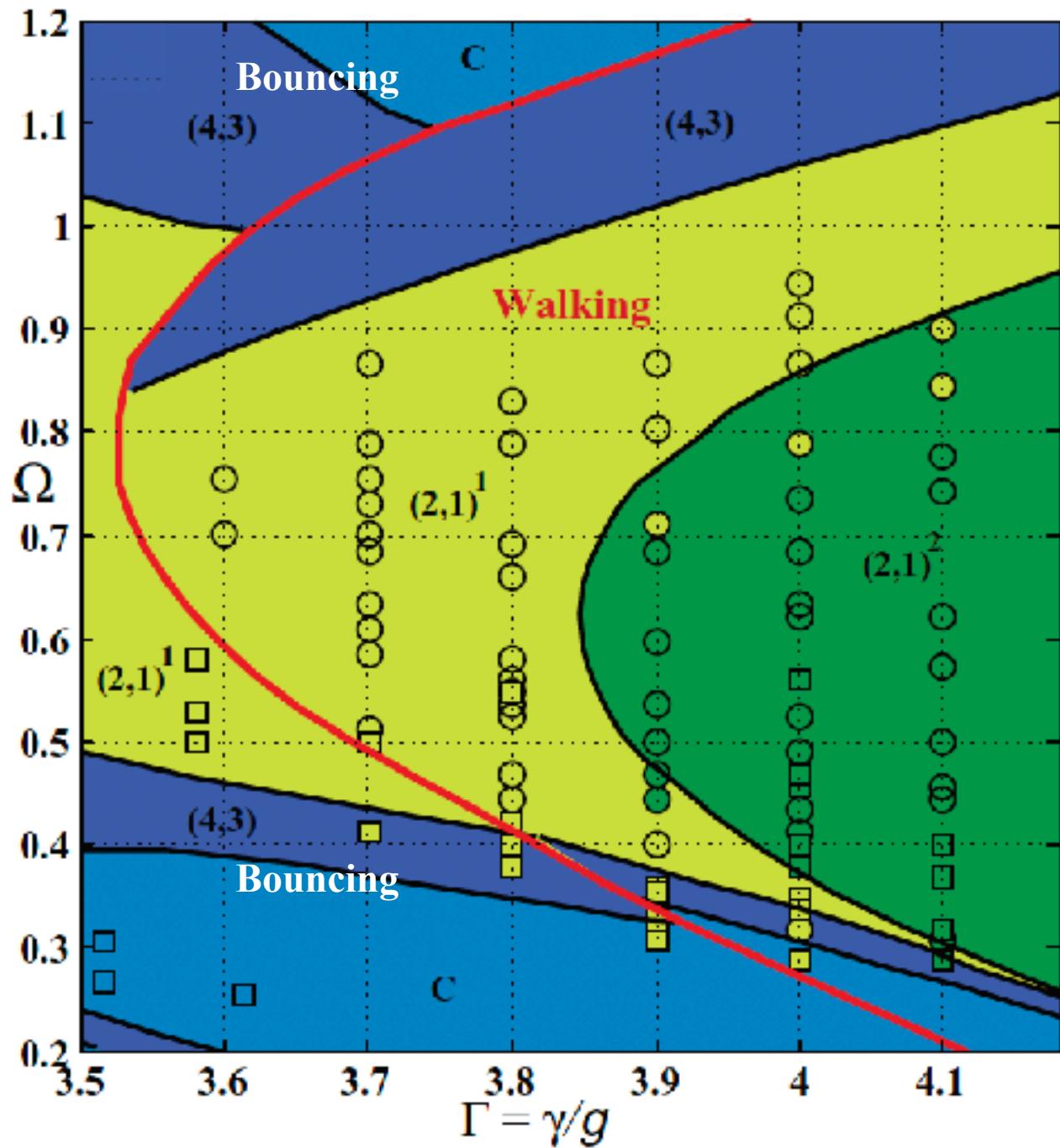
¹*Department of Applied Mathematics and Computer Science, Technical University
of Denmark, 2800 Kongens Lyngby, Denmark*

²*Department of Mathematics, Massachusetts Institute of Technology,
77 Massachusetts Avenue, Cambridge, Massachusetts 02139, USA*

Regime diagram

(50cS, 50Hz)

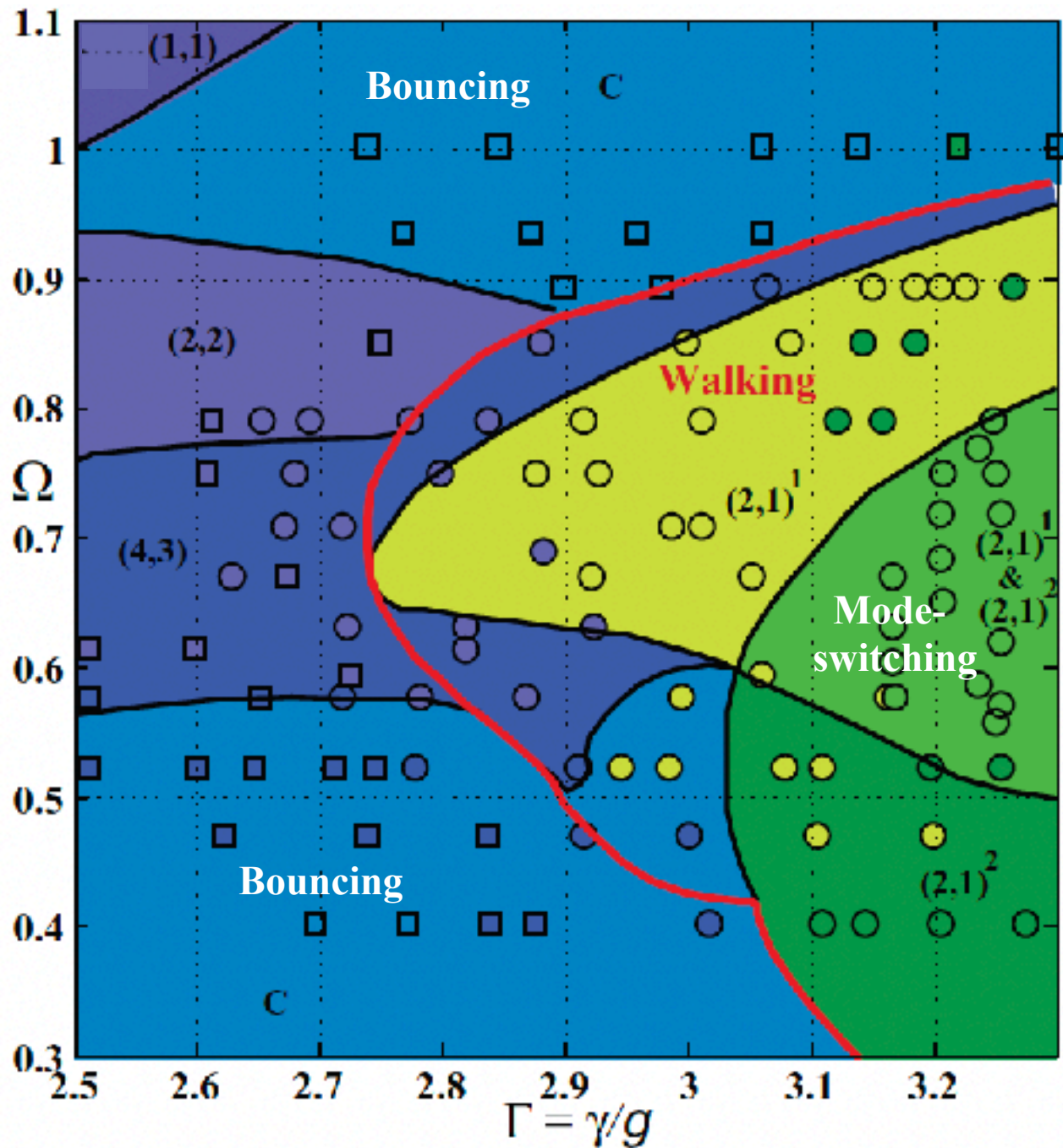
$$\Omega = 2\pi f(\rho R_0^3/\sigma)^{1/2}$$



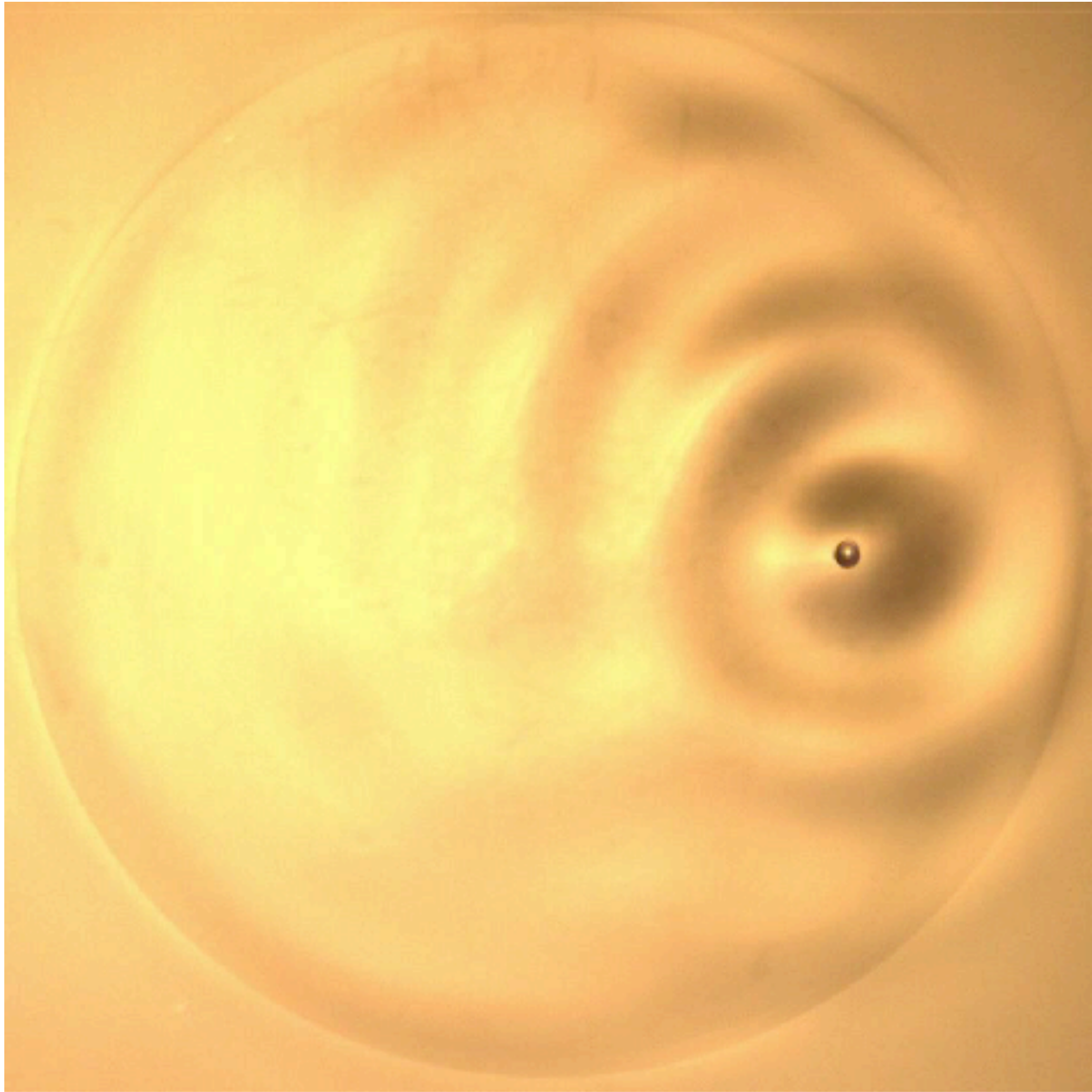
Regime diagram

(20cS, 70Hz)

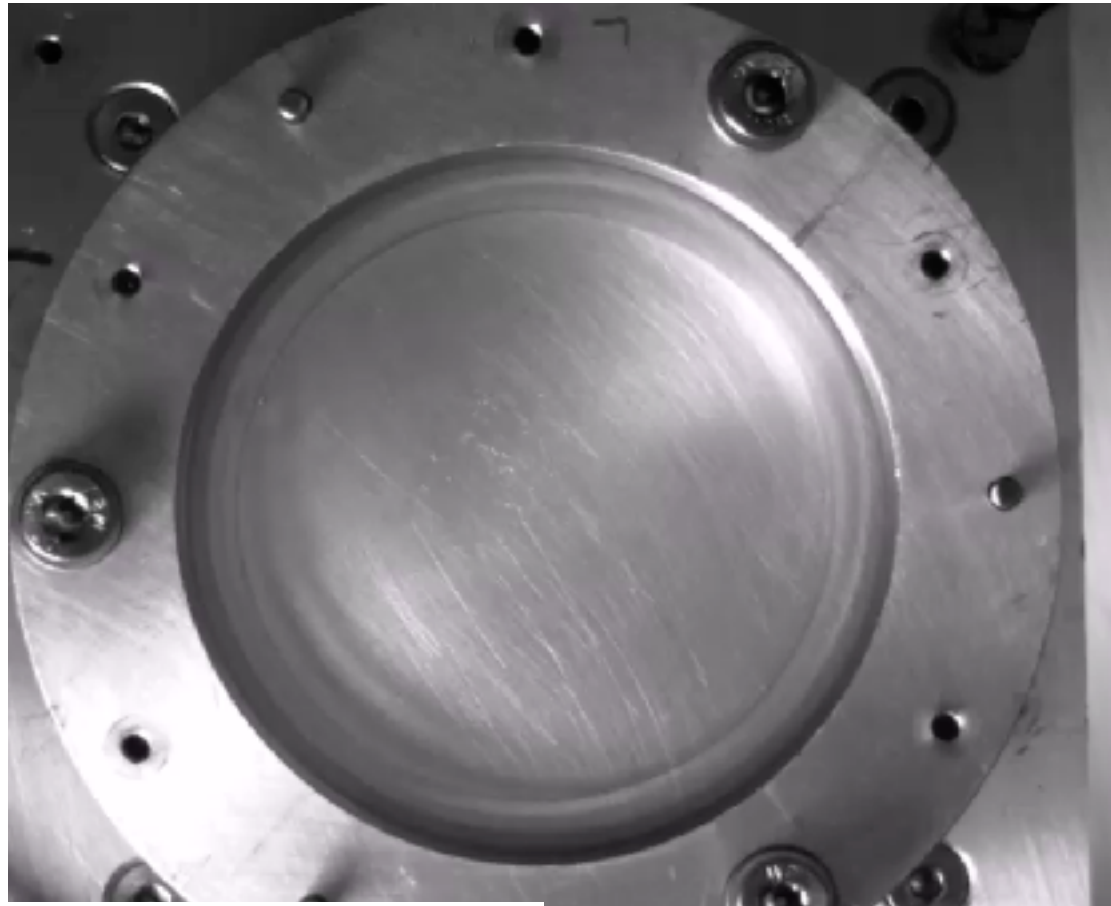
$$\Omega = 2\pi f(\rho R_0^3/\sigma)^{1/2}$$



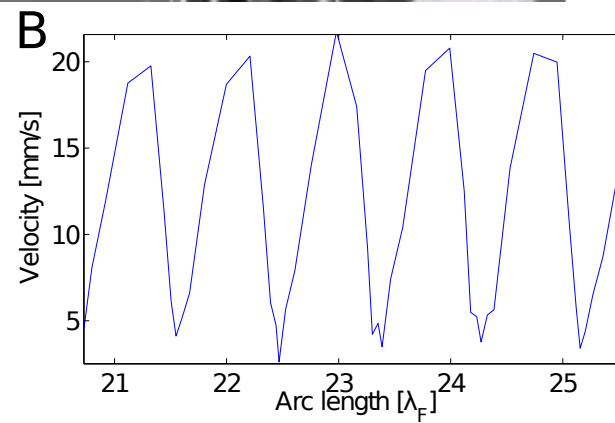
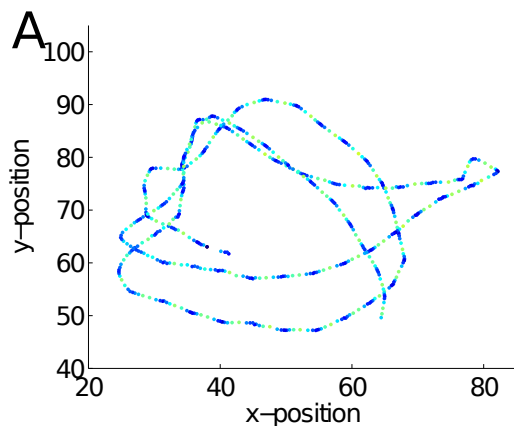
A walker in a mixed, mode-switching state



A walker in a mixed, mode-switching state



$(20cS, 70Hz)$



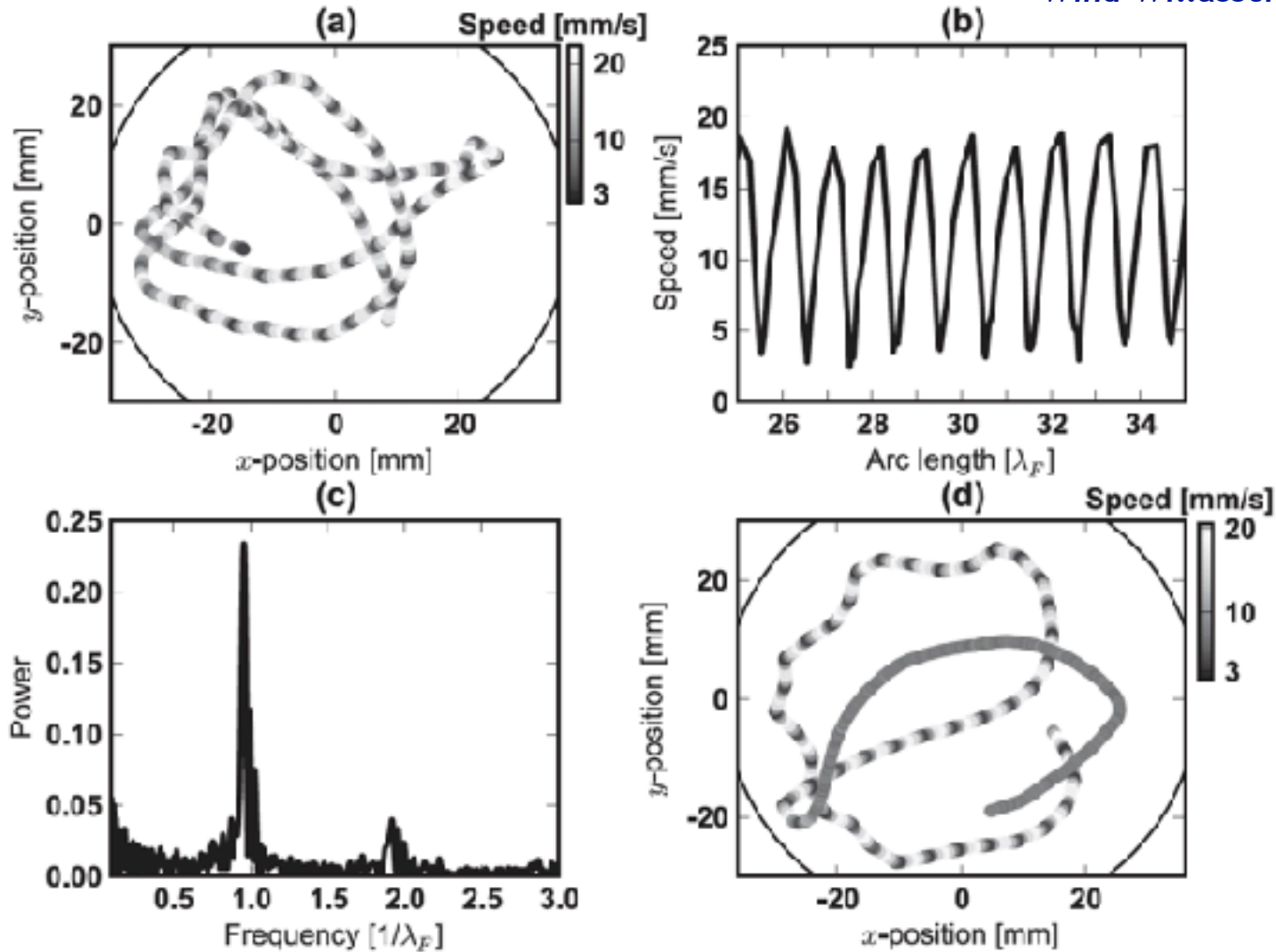
$(2,1)^2$

$(2,1)^1$

A walker in a mixed, mode-switching state

(20cS, 70Hz)

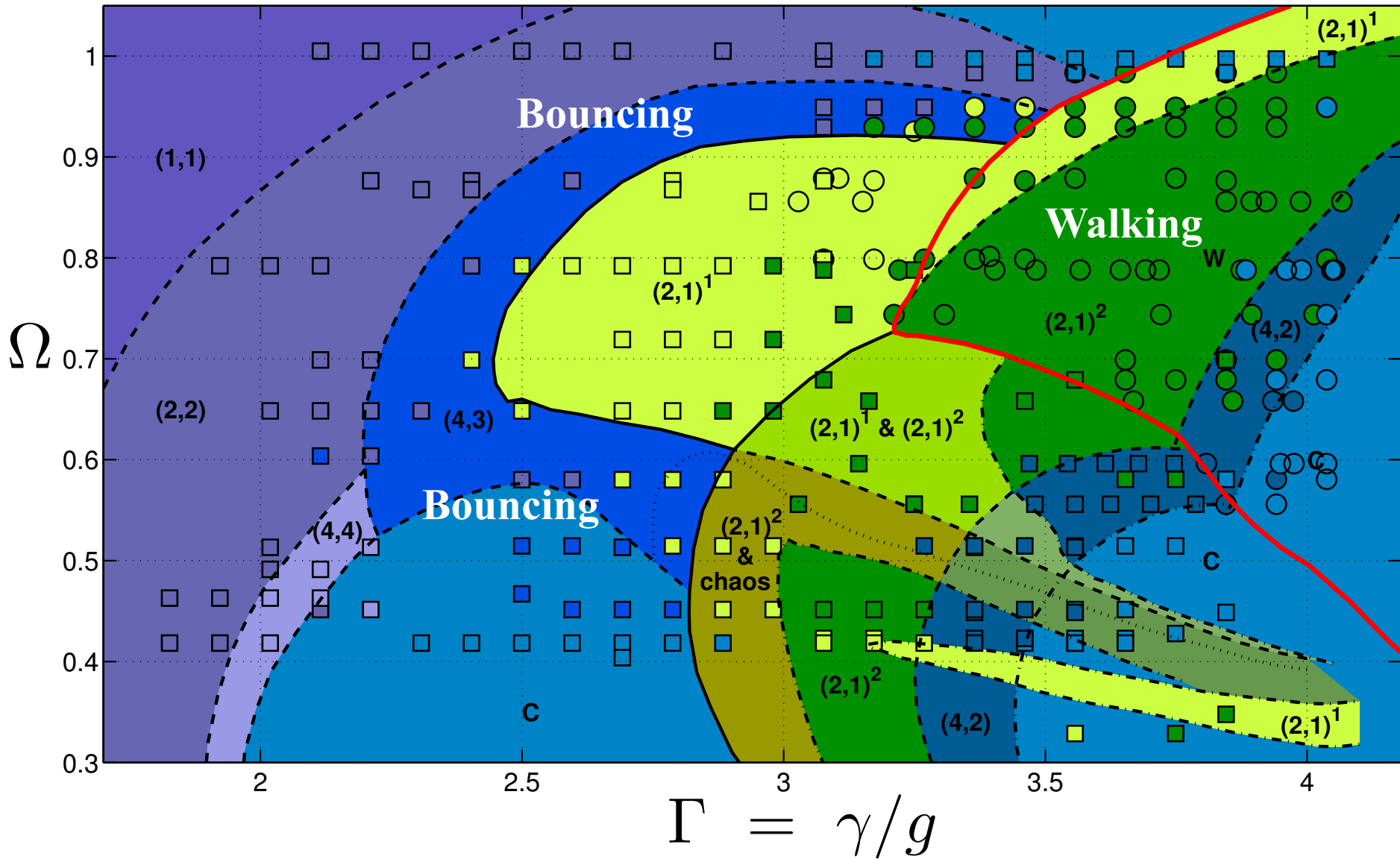
• *Wind-Willassen et al. (2013)*



- gives rise to in-line speed oscillations with the Faraday wavelength
- one of the three paradigms for the emergence of quantum-like statistics

Regime diagram (20cS, 80Hz)

$$\Omega = 2\pi f(\rho R_0^3/\sigma)^{1/2}$$



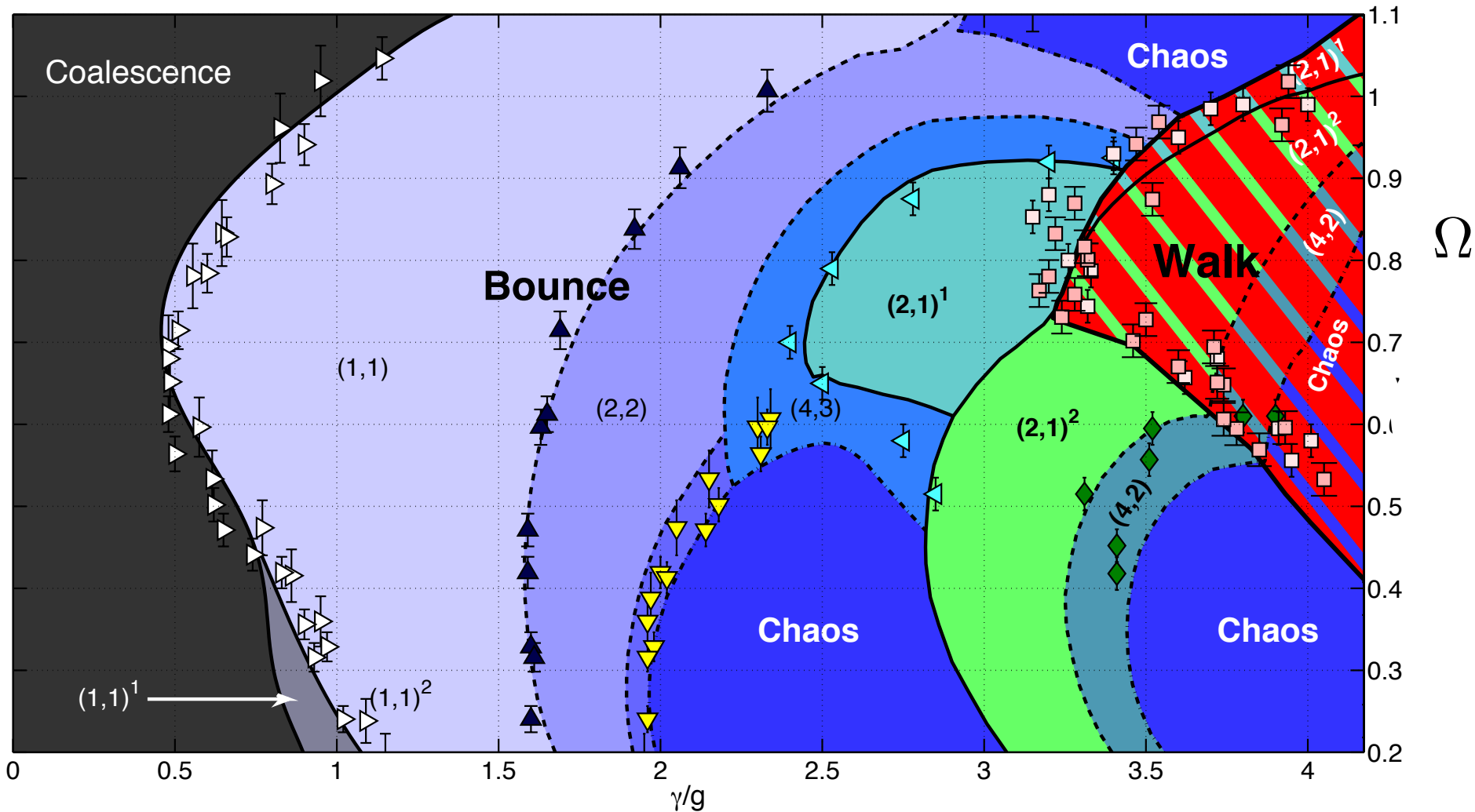
• Wind-Willassen, Molacek, Harris & Bush, *Phys. Fluids* (2013)

Molacek model of walking drops

(20cS, 80Hz)

Molacek & Bush (2013ab)

$$\Omega = 2\pi f(\rho R_0^3/\sigma)^{1/2}$$



In (m, n) mode, a drop bounces n times in m forcing periods.

- rationalized bouncing, walking behavior, but is cumbersome

Summary

In (the extremely limited) parameter regime of interest for walkers:

- drop deformation negligible
- interface behaves like a linear spring
- horizontal drag primarily due to momentum transferred to bath
- aerodynamics effects negligible in vertical dynamics, appreciable in horizontal
- drop impact generates Bessel function waves with spatiotemporal damping
- have rationalized regime diagrams, observed walking speeds
- formulation ungainly for non-resonant walkers, involving discrete sums
- time-averaging will yield the analytically tractable **stroboscopic model**

